

Chapter 4

Maxwell's Equations: The base of modern Electromagnetism is a set of four equations:

$$1) \vec{\nabla} \cdot \vec{D} = \rho_v$$

$$2) \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$3) \vec{\nabla} \cdot \vec{B} = 0$$

$$4) \vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

E : electric field

$$D = \epsilon E$$

B : magnetic field

$$B = \mu H$$

ρ_v : electric charge density per unit volume

J : current density

Published in 1873 by James Clerk Maxwell.

In static field all $\frac{\partial}{\partial t}$'s are zero:

Electrostatics:

$$\vec{\nabla} \cdot \vec{D} = \rho_v$$

$$\vec{\nabla} \times \vec{E} = 0$$

Magnetostatics:

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{H} = \vec{J}$$

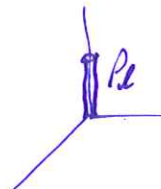
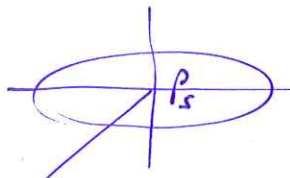
So in static case, E and B are no longer connected.

Charge and Current Distributions

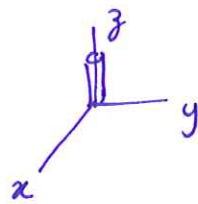
charge density: - volume charge density $\rho_v = \frac{dq}{dv} \left(\frac{C}{m^3}\right) \rightarrow Q = \int_V \rho_v dv \quad (C)$

- surface charge density $\rho_s = \frac{dq}{ds} \left(\frac{C}{m^2}\right) \rightarrow Q = \int_S \rho_s ds \quad (C)$

- line charge density $\rho_l = \frac{dq}{dl} \left(\frac{C}{m}\right) \rightarrow Q = \int_L \rho_l dl \quad (C)$



Ex: $\rho_l = 2z$
on a tube



if $l = 10$ cm, what is the total charge on this tube?

Answer: $Q = \int \rho_l dl = \int_0^{0.10} 2z dz = z^2 \Big|_0^{0.1} = 0.01$ C

Ex.

The charge on a disk is linearly increasing with r from zero at center to 6 C/m^2 at $r = 3$ cm. Find the total charge on this disk.

Answer:

$\rho_s = \frac{6r}{R} \rightarrow Q = \int \rho_s ds = \int \frac{6r}{R} r dr d\phi$

$Q = \int \frac{6r}{0.03} r dr d\phi = \int_{\phi=0}^{2\pi} d\phi \int_0^{0.03} 200r^2 dr = 2\pi \times 200 \times \frac{r^3}{3} \Big|_0^{0.03} = 11.31 \times 10^{-3}$ (C)

Current density:

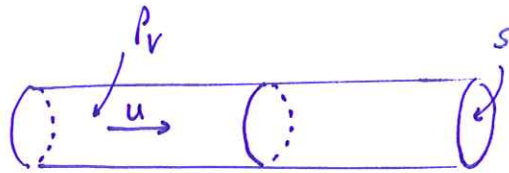
Current is the amount of charge moved per unit time:

$$I = \frac{dq}{dt}$$

current density is J and is current per unit area:

$$J = \frac{dI}{ds} \rightarrow I = \int_s \vec{J} \cdot d\vec{s}$$

We can also define current for the charges moving in a conductor with velocity u :



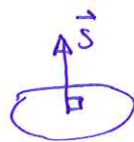
$$I = \frac{dq}{dt} = \frac{dq}{dv} \frac{dv}{dt} \rightarrow I = \rho_v \frac{sdl}{dt} = \rho_v u \rightarrow \frac{I}{s} = \rho_v u$$

$$\rho_v = \frac{dq}{dv}$$

$$\frac{dl}{dt} = u$$

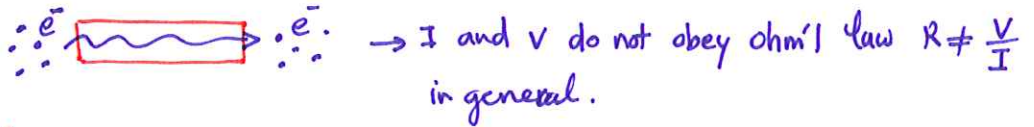
$$\vec{J} = \rho_v \vec{u} \quad \left(\frac{A}{m^2} \right)$$

Note that $d\vec{s}$ and \vec{s} are vectors normal to the surface:



* Convection Current:

Current is generated by actual movement of charges:



* Conduction Current:

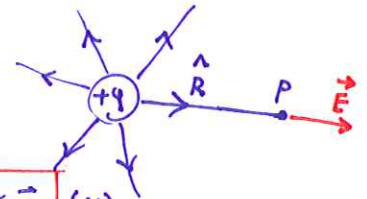
Electrons move within atoms and generate current. The electrons that emerge from the wire are not necessarily the same as the ones entered the wire at the other end:



Coulomb's Law

(1) \vec{E} for a charge q is given by:

$$\vec{E} = \hat{R} \frac{q}{4\pi\epsilon R^2}$$



(2) The force \vec{F} on a charge q in an Electric field \vec{E} is:

$$\vec{F} = q\vec{E} \quad (N)$$

\vec{E} : Electric field intensity

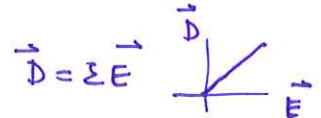
\vec{D} : Electric flux density

\vec{D} and \vec{E} are related with $\epsilon = \epsilon_r \epsilon_0$:

$$\vec{D} = \epsilon \vec{E}$$

where $\epsilon_0 = 8.85 \times 10^{-12} \approx \frac{10^{-9}}{36\pi} \quad (F/m)$

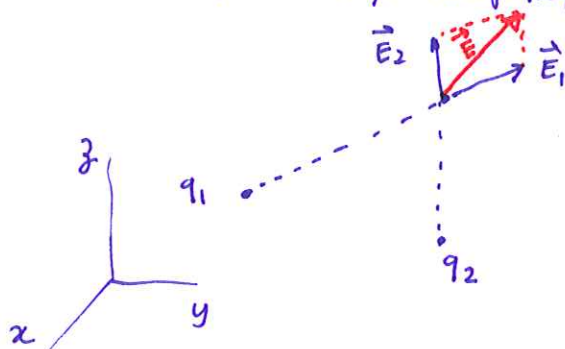
Linear Material: ϵ is independent of the magnitude of \vec{E} .



Isotropic Material: ϵ is independent of the direction of \vec{E} .

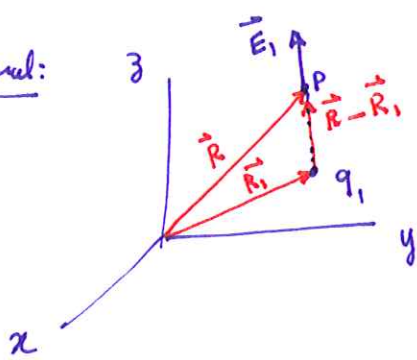
Electric Field due to Multiple Point Charges

The Electric field obeys the principle of linear superposition:



$$\vec{E} = \vec{E}_1 + \vec{E}_2$$

In general:



$$\vec{E}_1 = \frac{q_1}{4\pi\epsilon_0 |\vec{R} - \vec{R}_1|^2} \frac{\vec{R} - \vec{R}_1}{|\vec{R} - \vec{R}_1|}$$

$$\vec{E}_1 = \frac{q_1 (\vec{R} - \vec{R}_1)}{4\pi\epsilon_0 |\vec{R} - \vec{R}_1|^3} \left(\frac{V}{m}\right)$$

\vec{R} is the vector for point P.

\vec{R}_1 is the vector for the charge location (q_1 location)

So if we have two charges of q_1 at \vec{R}_1 and q_2 at \vec{R}_2 , the electric field at point \vec{R} is sum of \vec{E}_1 and \vec{E}_2 :

$$\vec{E} = \vec{E}_1 + \vec{E}_2 = \frac{q_1 (\vec{R} - \vec{R}_1)}{4\pi\epsilon_0 |\vec{R} - \vec{R}_1|^3} + \frac{q_2 (\vec{R} - \vec{R}_2)}{4\pi\epsilon_0 |\vec{R} - \vec{R}_2|^3}$$

And if we have N charges of q_1, q_2, \dots, q_N at locations $\vec{R}_1, \vec{R}_2, \dots, \vec{R}_N$, \vec{E} is:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^N \frac{q_i (\vec{R} - \vec{R}_i)}{|\vec{R} - \vec{R}_i|^3}$$

Example

$q_1 = 2 \times 10^{-5} \text{ C}$ is located at $(1, 3, -1)$ and $q_2 = -4 \times 10^{-5} \text{ C}$ at $(-3, 1, -2)$.

Find (a) \vec{E} at $(3, 1, -2)$ and (b) the force on a $8 \times 10^{-5} \text{ C}$ charge located at that point.

Answer:

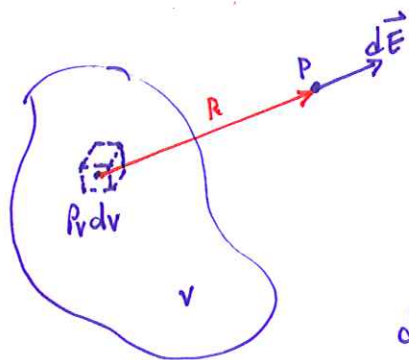
$$\begin{cases} \vec{R}_1 = \hat{x} + 3\hat{y} - \hat{z} \\ \vec{R}_2 = -3\hat{x} + \hat{y} - 2\hat{z} \\ \vec{R} = 3\hat{x} + \hat{y} - 2\hat{z} \end{cases}$$

(a)

$$\begin{aligned} \vec{E} &= \frac{1}{4\pi\epsilon_0} \left(q_1 \frac{(\vec{R} - \vec{R}_1)}{|\vec{R} - \vec{R}_1|^3} + q_2 \frac{(\vec{R} - \vec{R}_2)}{|\vec{R} - \vec{R}_2|^3} \right) \\ &= \frac{1}{4\pi\epsilon_0} \left(\frac{2(2\hat{x} - 2\hat{y} - \hat{z})}{(2^2 + 2^2 + 1)^{3/2}} + \frac{(-4)(6\hat{x})}{6^3} \right) \times 10^{-5} \\ &= \frac{10^{-5}}{108\pi\epsilon_0} (\hat{x} - 4\hat{y} - 2\hat{z}) \end{aligned}$$

$$(b) \vec{F} = q_3 \vec{E} = 8 \times 10^{-5} \times \frac{10^{-5}}{108\pi\epsilon_0} (\hat{x} - 4\hat{y} - 2\hat{z}) = \frac{10^{-10}}{27\pi\epsilon_0} (2\hat{x} - 8\hat{y} - 4\hat{z}) \text{ (N)}$$

Electric Field due to a charge Distribution



We want to find the electric field due to a volume of charge.

If the volume charge density is ρ_V , we can write:

$$d\vec{E} = \hat{R} \frac{\rho_V dv}{4\pi\epsilon R^2} = \hat{R} \frac{\rho_V dv}{4\pi\epsilon R^2}$$

$$\rightarrow \vec{E} = \int_V d\vec{E} = \frac{1}{4\pi\epsilon} \int_V \hat{R} \frac{\rho_V dv}{R^2}$$

Note that both R and \hat{R} change with position.

We can write \vec{E} for three cases of volumetric, surface, and line charges:

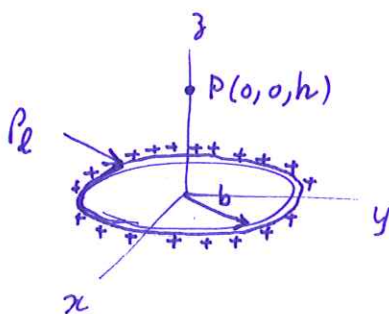
Volume:
$$\vec{E} = \frac{1}{4\pi\epsilon} \int_V \hat{R} \frac{\rho_V dv}{R^2}$$

Surface:
$$\vec{E} = \frac{1}{4\pi\epsilon} \int_S \hat{R} \frac{\rho_S ds}{R^2}$$

Line:
$$\vec{E} = \frac{1}{4\pi\epsilon} \int_L \hat{R} \frac{\rho_L dl}{R^2}$$

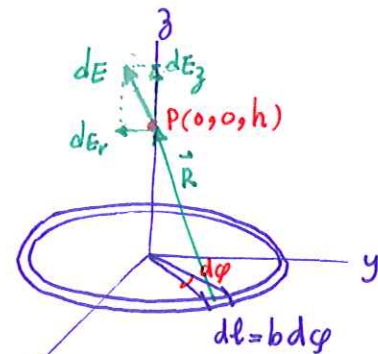
Example Electric field of a ring of charge:

Determine the Electric field for a ring of charge with uniform line charge density ρ_L at point $P(0,0,h)$ as shown in the picture:



Solution:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \hat{R} \frac{\rho_l dl}{R^2}$$



$$dl = b d\varphi$$

$$\hat{R} = \frac{-b\hat{r} + h\hat{z}}{\sqrt{b^2 + h^2}}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int (-b\hat{r} + h\hat{z}) \frac{\rho_l b d\varphi}{(b^2 + h^2)^{3/2}}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \left[\int_0^{2\pi} \frac{-b^2 \rho_l d\varphi}{(b^2 + h^2)^{3/2}} \hat{r} + \int_0^{2\pi} \frac{hb \rho_l d\varphi}{(b^2 + h^2)^{3/2}} \hat{z} \right]$$

= 0 from symmetry

$$= \frac{1}{4\pi\epsilon_0} \frac{hb\rho_l}{(b^2 + h^2)^{3/2}} \int_0^{2\pi} d\varphi \hat{z} = \frac{hb\rho_l}{2\epsilon_0(b^2 + h^2)^{3/2}} \hat{z}$$

Since the total charge is $Q = 2\pi b \rho_l \Rightarrow b\rho_l = \frac{Q}{2\pi}$ or $\vec{E} = \frac{hQ}{4\pi\epsilon_0(b^2 + h^2)^{3/2}} \hat{z}$

Example Electric Field of a circular disk of charge.

Find the electric field at a point $P(0,0,h)$ in free space due to a circular disk of charge in the x - y plane with uniform charge density ρ_s as shown in the picture:

For a ring of radius r and width dr we can write $d\vec{E}$ as:

$$d\vec{E} = \frac{h dq}{4\pi\epsilon_0(r^2 + h^2)^{3/2}} \hat{z} \quad dq = 2\pi r dr \rho_s$$

$$d\vec{E} = \frac{\rho_s h}{2\epsilon_0} \frac{r dr}{(r^2 + h^2)^{3/2}} \hat{z}$$

$$\vec{E} = \int d\vec{E} = \hat{z} \frac{\rho_s h}{2\epsilon_0} \int_0^a \frac{r dr}{(r^2 + h^2)^{3/2}} = \hat{z} \frac{\rho_s h}{4\epsilon_0} \int_{h^2}^{a^2 + h^2} \frac{du}{u^{3/2}} = \hat{z} \frac{\rho_s h}{4\epsilon_0} \left. \frac{-2}{\sqrt{u}} \right|_{h^2}^{a^2 + h^2}$$

($u = r^2 + h^2 \rightarrow du = 2r dr$)

$$\vec{E} = \hat{z} \frac{\rho_s h}{2\epsilon_0} \left(\frac{1}{\sqrt{h^2}} - \frac{1}{\sqrt{a^2 + h^2}} \right) = \hat{z} \frac{\rho_s}{2\epsilon_0} \left(\frac{h}{|h|} - \frac{h}{\sqrt{a^2 + h^2}} \right)$$

For infinite sheet of charge $a = \infty \Rightarrow \vec{E} = \hat{z} \frac{\rho_s}{2\epsilon_0}$ (For points with $z < 0$, \hat{z} must be replaced by $-\hat{z}$)

